# Buckling under the action of loading by aerodynamic and inertial forces during ground track tests of aviation equipment 

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#### Abstract

The article analyses the choice of a rational layout of the test object with a propulsion system (PS). One of the design examples of calculating the longitudinal stability and strength of the structure is given. The purpose of the article is to solve the problem of bending the elastic line of a cantilever tubular rod with a hinged termination during tests of a propulsion system for various aircrafts. On the example, the estimates of the approximate test object, accelerated on the track to a speed of $1200 \mathrm{~m} / \mathrm{s}$, are carried out. The aerodynamic loading of the structure of the mobile track installation is considered using the methods of mathematical modelling and the development of an algorithm for the numerical solution of the problem of bending the elastic line of a cantilever tubular rod. The deflection from the forces of external and internal loads of the outer shell of a movable track installation is considered, provided that the diameter of the outer contour is equal to the minimum and it is constant along the entire length.


Key Words: movable installation, construction, resistance forces

## 1. INTRODUCTION

The problems that arise before the designers of various types of aircraft require their consistent solution through ground track full-scale tests. To ensure reliability during high-speed testing of aircraft products, preliminary design estimates of the strength and stability of structural units are required, taking into account the influence of various factors when the maximum speed of the test object is reached. The practical implementation of ground track tests of aircraft objects is preceded by mathematical modelling and the development of an algorithm for the numerical solution of a problem that simulates test conditions [1], [2], [3]. The problems that arise before the designers of various types of aircraft require their consistent solution through ground track full-scale tests. To ensure reliability during high-speed testing of aircraft products, preliminary design estimates of the strength and stability of structural units are required, taking into account the influence of various factors when the maximum speed of the test object is reached. The practical implementation of ground track tests of
aircraft objects is preceded by mathematical modelling and the development of an algorithm for the numerical solution of a problem that simulates test conditions [4], [5]. Ground highspeed tests of various objects of aviation technology are carried out on an experimental track installation $3,500 \mathrm{~m}$ long. The track itself is a two-rail track that guides the movement of the test object placed on a track carriage, on two rails, or on a monorail [6], [7], [8], [9], [10].

The acceleration of the rail carriage is provided by a solid propellant rocket engine (SPRE). The test object is a hollow cylinder with a tapered fairing. The track carriage with welded-on shoes (conventionally sleeve bearings) can move freely along the rail. The shoes play the role of a hinged termination.

## 2. MATERIALS AND METHODS

Let us consider the aerodynamic loading of the structure of a mobile track installation. On the one hand, the thrust force of the SPRE is directed along the axis of the installation in the direction of travel; we denote it as $R$. On the other hand, the forces of aerodynamic resistance are directed against the direction of motion; we denote them as $P_{\Sigma}$. In this case, it is the total value of various types of aerodynamic drag, including wave drag. The forces of resistance to sliding friction caused by the normal force of pressure on the bearing surface of the shoes are also taken into account [6], [7], [11], [12], [13], as well as the loss of the traction force due to the creation of the lifting force due to the positive angle of inclination of the rail (angle of attack) when the carriage moves along the rail track [6], [14], [15]. The magnitude of the lifting force is variable during the movement of the rail carriage. The installation diagram is presented in Figure 1.


Fig. 1 - Image of the test object, placed on a track carriage with solid propellant motors
We consider a cantilever-mounted nose (a cylinder with a docked conical fairing without supports, in contrast to Fig. 1) and PS on two supports (shoes). One support is displaced 0.5 m from the joint with the nose, and the other retreats 0.5 m from the tail of the engine.

## 3. RESULTS AND DISCUSSIONS

Let us consider the deflection from the forces of external and internal loads of the outer shell of a movable track installation, provided that the diameter of the outer contour is equal to the minimum and is constant along the entire length. Equation of balance of forces:

$$
\begin{equation*}
P_{\text {Eaer }}(t)+R(t)=P_{\Sigma}(t) \tag{1}
\end{equation*}
$$

where: $R(t)$ is the thrust force of the $\operatorname{SPRE} ; P_{\text {Eaer }}(t)$ is total aerodynamic drag forces; $P_{\Sigma}(t)$ is the total compression force of the shell (tubular rod).

We denote the longitudinal axis of the mobile installation by $x$; it is directed parallel to the rail track along the movement.

The $y$-axis is perpendicular to the motion and is the direction of deflection of the elastic line of the tubular rod [16], [17], [18]. At small deformations within Hooke's law, the deflection equation will have the form [11]:

$$
\begin{equation*}
E J \frac{d^{2} y}{d x^{2}}= \pm M_{\mathrm{BEND}}(x) \tag{2}
\end{equation*}
$$

here: $E$ - modulus of elasticity; $J$ - the smallest moment of inertia of the section of the cylindrical part of the console. The bending moment in the current section $x$ will be equal to:

$$
\begin{equation*}
M_{\mathrm{BEND}}(x)=P_{\mathrm{LAT}} \cdot x-P_{\Sigma}(t) \cdot y \tag{3}
\end{equation*}
$$

$P_{\text {LAT }}$ is lateral force arising from the interaction of shock waves around the bow and reflected from the track surface when tested on a monorail (Fig. 2).


Fig. 2 - Image of the shock wave interference when the rail carriage moves at a speed of $1200 \mathrm{~m} / \mathrm{s}$. (The calculation was performed in Flow Vision)
Elastic line differential equation:

$$
\begin{equation*}
E J \frac{d^{2} y}{d x^{2}}=P_{\mathrm{LAT}} \cdot x-P_{\Sigma}(t) \cdot y \tag{4}
\end{equation*}
$$

The differential equation of the elastic line of the structure:

$$
\begin{equation*}
\frac{d^{2} y}{d x^{2}}+\frac{P_{\Sigma}(t)}{E J} \cdot y=\frac{P_{\mathrm{LAT}}}{E J} \cdot x \tag{5}
\end{equation*}
$$

The general solution to this differential equation is:

$$
\begin{equation*}
y=A \sin \sqrt{\frac{P_{\Sigma}(t)}{E J}} x+B \cos \sqrt{\frac{P_{\Sigma}(t)}{E J}} x+\frac{P_{\mathrm{LAT}}(t)}{P_{\Sigma}(t)} x \tag{6}
\end{equation*}
$$

for simplicity, we introduce the notation: the angle $\alpha$ is represented as follows:

$$
\begin{equation*}
\alpha=\sqrt{\frac{P_{\Sigma}(t)}{E J}} \tag{7}
\end{equation*}
$$

and the equation of the elastic line is written in the form:

$$
\begin{equation*}
y=A \sin \alpha x+B \cos \alpha x+\frac{P_{\mathrm{LAT}}}{P_{\Sigma}} \cdot x \tag{8}
\end{equation*}
$$

For an arbitrary moment in time, let us estimate the values of the constants [19]. Provided that the shoe touches the bottom surface of the rail head for the moment of time of reaching the maximum speed and on the ascending section of the track with an inclination angle equal to $\alpha=1.15^{\circ}(\tan \alpha=0.02)$, we assume that at $x=0$ the transverse deflection is zero, i.e., $y=$ 0 ; from here follows $B=0$. For the case under consideration, shoes No. 1, No. 2, No. 4 (see Fig. 1) do not exist. We consider shoe No. 3 to be a hinged termination. Bending of the bar is realised along the curve of least resistance (least stiffness). Interest in the problem under consideration was manifested in works [20], [21], [22], [23], [24], [25], [26].

Next, we consider the problem of bending the elastic line of a cantilever tubular rod bounded by support No. 3, which we now consider to be a hinged termination. The elastic line equation [27] is described by the equation:

$$
\begin{equation*}
y=A \sin \alpha x+\frac{P_{\mathrm{LAT}}}{P_{\Sigma}} \cdot x \tag{9}
\end{equation*}
$$

An approximate deflection can be represented by the dependence (Fig. 3):

$$
\begin{equation*}
y(x)=a_{0}+a_{1} x+a_{2} x^{2}+a_{3} x^{3}+\ldots+a_{n} x^{n} \tag{10}
\end{equation*}
$$



Fig. 3 - Simplified diagram of the deflection of an elastic line.
To determine the coefficients, we formulate the boundary conditions for the considered scheme: $y(0)=0 ; y^{\prime \prime}(0)=0 ; \mathrm{y}\left(l_{2}\right)=0 ; \mathrm{y}^{\prime \prime}\left(l_{1}+l_{2}\right)=0 ; y\left(l_{1}+l_{2}\right)=\Delta$ is the required deflection of the fairing apex from the axis. From the first two conditions it follows that $a_{0}=0, a_{2}=0$. From the third and fourth conditions we obtain:

$$
\begin{gather*}
y\left(l_{2}\right)=a_{1} l_{2}+a_{3} l_{2}^{3}+a_{4} l_{2}^{4}=0  \tag{11}\\
y^{\prime \prime}(x)=a_{3}+2 a_{4}\left(l_{1}+l_{2}\right)=0 \tag{12}
\end{gather*}
$$

Let us express the coefficients:

$$
\begin{gather*}
a_{3}=-2 a_{4}\left(l_{1}+l_{2}\right),  \tag{13}\\
a_{1} l_{2}+\left(-2 a_{4}\left(l_{1}+l_{2}\right)\right) l_{2}^{3}+a_{4} l_{2}^{4}=0,  \tag{14}\\
a_{1}=a_{4}\left[2\left(l_{1}+l_{2}\right) l_{2}^{2}-l_{2}^{3}\right]=a_{4}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right) \tag{15}
\end{gather*}
$$

The equation of the elastic line of the mobile unit will take the form:

$$
\begin{gather*}
y(x)=a_{4}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right) x-2 a_{4}\left(l_{1}+l_{2}\right) x^{3}+a_{4} x^{4}  \tag{16}\\
y^{\prime}(x)=a_{4}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right)-6 a_{4}\left(l_{1}+l_{2}\right) x^{2}+4 a_{4} x^{3}  \tag{17}\\
y^{\prime \prime}(x)=12 a_{4} x^{2}-12 a_{4}\left(l_{1}+l_{2}\right) x=12 a_{4}\left[x^{2}-\left(l_{1}+l_{2}\right) x\right]  \tag{18}\\
{\left[y^{\prime}(x)\right]^{2}=a_{4}^{2}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right)^{2}+36 a_{4}^{2}\left(l_{1}+l_{2}\right)^{2} x^{4}+16 a_{4}^{2} x^{6}-}  \tag{19}\\
-12 a_{4}^{2}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right)\left(l_{1}+l_{2}\right)+8 a_{4}^{2}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right) x^{3}-48 a_{4}^{2}\left(l_{1}+l_{2}\right) x^{3}
\end{gather*}
$$

$$
\begin{gather*}
{\left[y^{\prime}(x)=16 a_{4}^{2} x^{6}-48 a_{4}^{2}\left(l_{1}+l_{2}\right) x^{5}+36 a_{4}^{2}\left(l_{1}+l_{2}\right)^{2} x^{4}+8 a_{4}^{2}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right) x^{3}\right.} \\
-12 a_{4}^{2}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right)\left(l_{1}+l_{2}\right) x^{2}+a_{4}^{2}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right)^{2}  \tag{20}\\
{\left[y^{\prime \prime}(x)\right]^{2}=144 a_{4}^{2}\left[x^{4}-2\left(l_{1}+l_{2}\right) x^{3}+\left(l_{1}+l_{2}\right)^{2} x^{2}\right],} \tag{21}
\end{gather*}
$$

$$
\begin{equation*}
\int_{0}^{l} E J\left(y^{\prime \prime}\right)^{2} d x=E J 144 a_{4}^{2} \int_{0}^{l}\left[x^{4}-2\left(l_{1}+l_{2}\right) x^{3}+\left(l_{1}+l_{2}\right)^{2} x^{2}\right] d x \tag{22}
\end{equation*}
$$

$$
\begin{equation*}
\int_{0}^{l} E J\left(y^{\prime \prime}\right)^{2} d x=E J 144 a_{4}^{2}\left[\frac{l^{5}}{5}-\frac{2\left(l_{1}+l_{2}\right) l^{4}}{4}+\frac{\left(l_{1}+l_{2}\right)^{2} l^{3}}{3}\right] \tag{23}
\end{equation*}
$$

$$
\begin{equation*}
\int_{0}^{l}\left(y^{\prime}\right)^{2} d x=a_{4}^{2}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right)^{2} \tag{24}
\end{equation*}
$$

$$
\int_{0}^{l}\left(y^{\prime}\right)^{2} d x=a_{4}^{2}\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right)^{2}+
$$

$$
\begin{equation*}
+a_{4}^{2}\left\{\frac{16 x^{7}}{7}-\frac{48\left(l_{1}+l_{2}\right) l^{6}}{6}+\frac{36\left(l_{1}+l_{2}\right)^{2} l^{5}}{5}+\frac{8\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right) l^{4}}{4}\right. \tag{25}
\end{equation*}
$$

$$
\left.-\frac{12\left(2 l_{1} \cdot l_{2}^{2}+l_{2}^{3}\right)\left(l_{1}+l_{2}\right) l^{3}}{3}\right\}
$$

The critical force is determined from the expression [11]:

$$
\begin{equation*}
P_{\mathrm{n}_{-} \mathrm{cr}}=\int_{0}^{l} E J\left(y^{\prime \prime}\right)^{2} d x / \int_{0}^{l}\left(y^{\prime}\right)^{2} d x \tag{26}
\end{equation*}
$$

It is equal to:

$$
\begin{equation*}
P_{\mathrm{n}-\mathrm{cr}}=\frac{144 E]\left[\left[_{5}^{5}-\frac{2\left(l_{1}+l_{2}\right) l^{4}}{4}+\frac{\left(l_{1}+l_{2}\right)^{2} l^{3}}{3}\right]\right.}{\left(2 l_{1} \cdot l_{2}{ }^{2}+l_{2}{ }^{3}\right)^{2}+\left\{\frac{16 l^{7}}{7}-\frac{48\left(l_{1}+l_{2}\right) l^{6}}{6}+\frac{36\left(l_{1}+l_{2}\right)^{2} l^{5}}{5}+\frac{8\left(2 l_{1} \cdot l_{2}{ }^{2}+l_{2}^{3}\right) l^{4}}{4}-\frac{12\left(2 l_{1} \cdot l_{2}{ }^{2}+l_{2}{ }^{3}\right)\left(l_{1}+l_{2}\right) l^{3}}{3}\right\}} \tag{27}
\end{equation*}
$$

where: $l_{1}=2260 \mathrm{~mm} ; l_{2}=4630 \mathrm{~mm}, P_{\mathrm{n}_{-} \text {cr }}=0,06 \cdot E J \frac{1}{\mathrm{~m}^{2}}$. Dimensionless length reduction factor is:

$$
\begin{equation*}
\mu=\pi \sqrt{\frac{\left(2 l_{1} \cdot l_{2}{ }^{2}+l_{2}^{3}\right)^{2}+\left\{\frac{16 l^{7}}{7}-\frac{48\left(l_{1}+l_{2}\right) l^{6}}{6}+\frac{36\left(l_{1}+l_{2}\right)^{2} l^{5}}{5}+\frac{8\left(2 l_{1} \cdot l_{2}{ }^{2}+l_{2}^{3}\right) l^{4}}{4}-\frac{12\left(2 l_{1} \cdot l_{2}{ }^{2}+l_{2}^{3}\right)\left(l_{1}+l_{2}\right) l^{3}}{3}\right\}}{144 \cdot l^{[ }\left[\frac{l^{5}}{5}-\frac{2\left(l_{1}+l_{2}\right) l^{4}}{4}+\frac{\left(l_{1}+l_{2}\right)^{2 l^{3}}}{3}\right]}}, \tag{28}
\end{equation*}
$$

i.e., $\mu=1,86$. For the aluminium alloy D16T [4], [5], [28], $E=0,75 \times 10^{5} \mathrm{MP}, \sigma_{\mathrm{pc}}=200 \mathrm{MPa}$. Moment of inertia of an unreinforced shell:

$$
\begin{align*}
& J_{x}=J_{y}=\frac{\pi D^{4}}{64}\left(1-C^{4}\right)  \tag{29}\\
& C=\frac{d}{D}=\frac{274}{280}=0,97857 \tag{30}
\end{align*}
$$

The moment of inertia at a shell thickness of 3 mm is: $J_{x}=J_{y}=2504.264 \mathrm{~cm}^{4}, E J_{x}=$ $E J_{y}=1,878198 \cdot 10^{6} \mathrm{H} \cdot \mathrm{m}^{2}$. The cross-sectional area of the shell $F=0.0026107 \mathrm{~m}^{2}$.

The smallest radius of gyration of the cross-section of a thin-walled shell [29] is: $i=$ $\sqrt{\frac{J}{F}}=0.09794 m$. Element flexibility:

$$
\begin{equation*}
\lambda=\frac{\mu l}{i} \tag{31}
\end{equation*}
$$

For the entire installation:

$$
\begin{equation*}
\lambda=\frac{\mu l}{i}=\frac{1.86 \cdot 6.89}{0.09794}=\frac{12.8154}{0.09794}=130.85 \tag{32}
\end{equation*}
$$

Body flexibility $\lambda_{0}=60, \lambda>\lambda_{0}$. Euler's approach is applicable [11] for the forebody:

$$
\begin{align*}
\lambda & =\frac{\mu l}{i}=\frac{1.86 \cdot 2.26}{0.09794}=\frac{4.2036}{0.09794}=42.9  \tag{33}\\
P_{\mathrm{n}_{\mathrm{c}} \mathrm{cr}} & =0,06 \cdot 1,878198 \cdot 10^{6}=112.692 \mathrm{kN}, \tag{34}
\end{align*}
$$

The displacement of the tip of the cone $\Delta$ caused by the bending of the elastic axis is determined approximately from the triangle with the hypotenuse $d x$ [11]. The displacement $\delta l$ of the apex of the fairing cone along the $x$-axis due to the inclination of the elastic axis is equal to:

$$
\begin{equation*}
\delta l=d x-d x \cos \alpha=2 \sin ^{2}(\alpha / 2) d x \tag{35}
\end{equation*}
$$

or

$$
\begin{equation*}
\delta l=\frac{\alpha^{2}}{2}=\frac{\left(y^{\prime}\right)^{2}}{2} d x \tag{36}
\end{equation*}
$$

Summing $\delta l$ along the length of the rod, we obtain the displacement along the $x$-axis due to bending:

$$
\begin{gather*}
\delta=\frac{1}{2} \int_{0}^{l}\left(y^{\prime}\right)^{2} d x,  \tag{37}\\
\alpha=\sqrt{\frac{P_{\mathrm{n}_{-} \mathrm{cr}}}{E J}}=0.244949 \frac{1}{m} . \tag{38}
\end{gather*}
$$

The curvilinear shape of the cantilever corresponds to $\cos \alpha \cdot l=0$ or $\alpha \cdot l_{1}=\pi / 2$ [30]. When the maximum speed is reached, the total load is equal to $P_{\Sigma}(t)=536 \mathrm{kN}$. The critical force for the cantilever part $n=2$ is

$$
\begin{equation*}
P_{\mathrm{n}_{-} \mathrm{cr} 1}=\frac{\pi^{2} E J}{n\left(\mu l_{1}\right)^{2}}=1503.3 k N \tag{39}
\end{equation*}
$$

and the whole tubular rod:

$$
\begin{equation*}
P_{\text {cr } \max }=\frac{\pi^{2} E J}{\left(\mu\left(l_{1}+l_{2}\right)\right)^{2}}=712 \mathrm{kN} \tag{40}
\end{equation*}
$$

The stability condition $P_{\Sigma}(t)<P_{\text {cr max }}$ is satisfied [31]. In this case, the stability margin is $n=1.33$. The bend of the cantilever bow corresponds to the dependence:

$$
\begin{equation*}
y=A \sin \frac{\pi x}{2 l_{1}} \tag{41}
\end{equation*}
$$

The displacement $\Delta$ along the y-axis can reach limiting values equal to $\Delta=6,7 \mathrm{~mm}$.

## 4. CONCLUSIONS

Based on the results of the study, using mathematical modelling, the problem of bending the elastic line of a cantilever tubular rod, limited by support No. 3, was solved. The deflection from the forces of external and internal loads of the outer shell of the movable track installation was also considered. The obtained calculations can be useful for strength testing of rocket launchers, as well as other aviation equipment of various types.

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